

For the details, see [JP],[JOPP],[AJPP] and the references therein. We consider the following setting of nonequilibrium statistical mechanics. The system we consider is given by a  $C^*$ -algebra

$$\mathcal{O} := \otimes_k \mathcal{O}_k,$$

associated with free time evolution

$$\tau_0^t = \otimes_k \tau_k^t.$$

The initial state is

$$\omega = \otimes_k \omega_k,$$

where each  $\omega_k$  is  $(\beta_k, \tau_k)$ -KMS state. The dynamics  $\tau$  is given by perturbing the free dynamics by an interaction  $V \in \mathcal{O}$ . If there exists a state  $\omega_+$  such that

$$\omega_+(A) = \lim_{t \rightarrow \infty} \omega_t(A), \quad A \in \mathcal{O},$$

we call it nonequilibrium steady state. We furthermore assume the existence of time reversal symmetry  $\Theta$ , i.e., an antilinear automorphism such that

$$\tau_t \circ \Theta = \Theta \circ \tau_{-t}, \quad \Theta^2 = id.$$

For a state  $\phi$ , its time reversal  $\bar{\phi}$  is given by  $\bar{\phi}(a) = \phi(\Theta(a^*))$ . We consider the situation that  $\bar{\omega} = \omega$ . We have  $\bar{\omega}_t = \omega_{-t}$ , so  $\omega_t$  is not time reversal invariant anymore, in general.

We want to estimate *how far*  $\omega_t$  and  $\omega_{-t}$  are. To do that, we consider measures like

$$D(\omega_1, \omega_2) := 1 - \frac{1}{2} \|\omega_1 - \omega_2\|,$$

$$\inf_{0 \leq T \leq 1} \{\omega_1(T) : \omega_2(T) \geq 1 - \varepsilon, \}, \quad 0 < \varepsilon < 1.$$

For  $s \in \mathbb{R}$ , we define

$$E_s(\omega_t, \omega_{-t}) = \log(\Delta_{\omega_t, \omega_{-t}}^s \xi_{\omega_{-t}}, \xi_{\omega_{-t}}).$$

Here  $\xi_\phi$  is the representative of  $\phi$  in the natural positive cone. The main result in [JOPS] is the following.

**Theorem 0.1** [JOPS] *Assume that the limit*

$$e(s) = \lim_{t \rightarrow \infty} \frac{1}{2t} E_s(\omega_t, \omega_{-t})$$

*exists and differentiable in some open interval  $I$  which includes  $[0, 1]$ , and  $e'(1) > 0$ . Then*

1.

$$\lim_{t \rightarrow \infty} \frac{1}{2t} \log D(\omega_t, \omega_{-t}) = e\left(\frac{1}{2}\right).$$

2.

$$\underline{B}(r) = \bar{B}(r) = B(r) = - \sup_{s \in [0, 1[} \frac{-sr - e(s)}{1 - s}.$$

for

$$\bar{B}(r) = \inf_{\{T_t\}} \left\{ \limsup_{t \rightarrow \infty} \frac{1}{2t} \log \omega_t(1 - T_t) \mid \limsup_{t \rightarrow \infty} \frac{1}{2t} \log \omega_{(-t)}(T_t) < -r \right\},$$

$$\underline{B}(r) = \inf_{\{T_t\}} \left\{ \liminf_{t \rightarrow \infty} \frac{1}{2t} \log \omega_t(1 - T_t) \mid \limsup_{t \rightarrow \infty} \frac{1}{2t} \log \omega_{(-t)}(T_t) < -r \right\},$$

$$B(r) = \inf_{\{T_t\}} \left\{ \lim_{t \rightarrow \infty} \frac{1}{2t} \log \omega_t(1 - T_t) \mid \limsup_{t \rightarrow \infty} \frac{1}{2t} \log \omega_{(-t)}(T_t) < -r \right\}.$$

where in the last case the infimum is taken over all families of tests  $\{T_t\}$  for which

$$\lim_{t \rightarrow \infty} \frac{1}{2t} \log \omega_t(1 - T_t)$$

exists.

3. For  $\epsilon \in ]0, 1[$ ,

$$\bar{B}_\epsilon = \underline{B}_\epsilon = B_\epsilon = -e'(1).$$

where

$$\bar{B}_\epsilon = \inf_{\{T_t\}} \left\{ \limsup_{t \rightarrow \infty} \frac{1}{2t} \log \omega_t(1 - T_t) \mid \omega_{(-t)}(T_t) \leq \epsilon \right\},$$

$$\underline{B}_\epsilon = \inf_{\{T_t\}} \left\{ \liminf_{t \rightarrow \infty} \frac{1}{2t} \log \omega_t(1 - T_t) \mid \omega_{(-t)}(T_t) \leq \epsilon \right\},$$

$$B_\epsilon = \inf_{\{T_t\}} \left\{ \lim_{t \rightarrow \infty} \frac{1}{2t} \log \omega_t(1 - T_t) \mid \omega_{(-t)}(T_t) \leq \epsilon \right\},$$

The rate  $e'(1)$  is called entropy production.

## References

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- [JP] Jaksic V., Pillet C.-A.: Mathematical theory of non-equilibrium quantum statistical mechanics, J. Stat. Phys. 108 (2002), 787-829
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