

# Conforming Vector Interpolation Functions for Polyhedral Meshes

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joint work with

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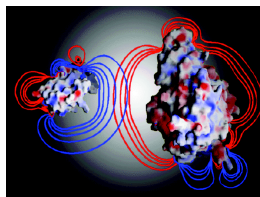
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# Interpolation in Graphics vs. Simulation



- Interpolation of vector fields required for geometric design.
- No natural constraints on interpolation properties.
- Some exploration of scalar interpolation over polyhedra.

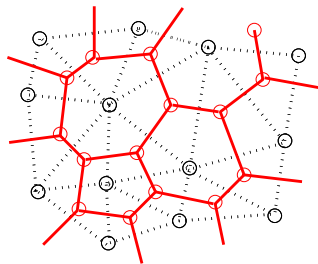
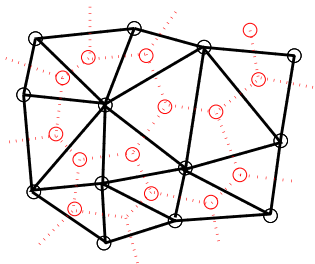


- Coupled vector fields related by integral and differential equations.
- Physical nature of problem offers natural discretizations of variables and boundary conditions.
- Discrete Exterior Calculus suggests a need for vector interpolation over polyhedra.

**Goal:** Develop a theory of vector interpolation over polyhedra conforming to physical requirements with provable error estimates.

# Motivation

Many authors (Bossavit, Hiptmair, Shashkov, . . .) have recognized the natural interplay between **primal** and **dual** domain meshes for discretization of physical equations.



Potential benefits of a theory based on interpolation over dual meshes:

- 1 Accuracy vs. speed tradeoffs available between primal and dual methods.
- 2 Error estimates for dual interpolation methods analogous to standard estimates.
- 3 Validation of primal-based results with dual-based discretization methods.

# Outline

- 1 Background on Vector Interpolation
- 2 Novel Discretizations Using Polyhedral Vector Interpolation
- 3 Error Estimates for Polyhedral Vector Interpolation

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# $H(\text{Curl})$ versus $H(\text{Div})$

Throughout, we will consider a model problem from **magnetostatics**:

- **Domain:** Contractible 3-manifold  $\Omega \subset \mathbb{R}^3$  with boundary  $\Gamma$

- **Variables:**

  - $b$  (magnetic field / magnetic induction)

  - $h$  (magnetizing field / auxiliary magnetic field)

- **Input:**

  - $j$  (current density field)

- **Equations:**

$$\operatorname{div} b = 0, \quad *b = h, \quad \operatorname{curl} h = j$$

- **Boundary Conditions:**  $\Gamma$  written as a disjoint union  $\Gamma^e \cup \Gamma^h$  such that

$$\hat{n} \cdot b = 0 \text{ on } \Gamma^e, \quad \hat{n} \times h = 0 \text{ on } \Gamma^h.$$

While  $b$  and  $h$  are both discretized as vector fields, they lie in different function spaces:

$$h \in H(\operatorname{curl}) := \left\{ \vec{v} \in \left( L^2(\Omega) \right)^3 \text{ s.t. } \nabla \times \vec{v} \in \left( L^2(\Omega) \right)^3 \right\}$$

$$b \in H(\operatorname{div}) := \left\{ \vec{v} \in \left( L^2(\Omega) \right)^3 \text{ s.t. } \nabla \cdot \vec{v} \in L^2(\Omega) \right\}$$

# Local Conformity Constraints

- Functional continuity can be enforced on a mesh  $\mathcal{T}$  by imposing certain constraints at each face  $F = T_1 \cap T_2$ , involving the normals to the mesh elements  $T_1, T_2$ :

$$H(\text{curl}) := \left\{ \vec{v} \in \left( L^2(\Omega) \right)^3 \quad \text{s.t.} \quad \nabla \times \vec{v} \in \left( L^2(\Omega) \right)^3 \right\}$$
$$h \in H(\text{curl}) \iff h|_{T_1} \times \hat{n}_1 + h|_{T_2} \times \hat{n}_2 = 0, \quad \forall F \in \mathcal{T}$$

$$H(\text{div}) := \left\{ \vec{v} \in \left( L^2(\Omega) \right)^3 \quad \text{s.t.} \quad \nabla \cdot \vec{v} \in L^2(\Omega) \right\}$$
$$b \in H(\text{div}) \iff b|_{T_1} \cdot \hat{n}_1 + b|_{T_2} \cdot \hat{n}_2 = 0, \quad \forall F \in \mathcal{T}$$

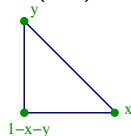
- These constraints hold for primal meshes ( $T_i$ =tetrahedra) **and** dual meshes ( $T_i$ =polyhedra).
- Goal:** Solve for  $h$  and  $b$  as functions defined piecewise over  $\mathcal{T}$ , guaranteed to satisfy the applicable conformity constraints.

# Whitney Elements for Primal Meshes

- The **Whitney elements** provide a simple and canonical way to construct piecewise functions over a **primal** mesh  $\mathcal{T}$  in  $H(\text{curl})$  or  $H(\text{div})$ :

- Start with linear barycentric coordinates:

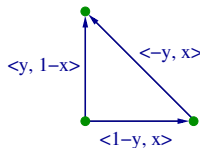
$$\lambda_i(\mathbf{v}_j) = \delta_{ij}$$



$\lambda_i \rightarrow 1$  d.o.f. per vertex

- Define for each edge  $\mathbf{v}_i\mathbf{v}_j$ :

$$\eta_{ij} := \lambda_i \nabla \lambda_j - \lambda_j \nabla \lambda_i$$



$\eta_{ij} \rightarrow 1$  d.o.f. per edge

- Define for each face  $\mathbf{v}_i\mathbf{v}_j\mathbf{v}_k$ :

$$\eta_{ijk} := \lambda_i \nabla \lambda_j \times \nabla \lambda_k + \lambda_j \nabla \lambda_k \times \nabla \lambda_i$$

$$+ \lambda_k \nabla \lambda_i \times \nabla \lambda_j$$

$$\langle -x, -y, 1-z \rangle, \langle x, y-1, z \rangle$$

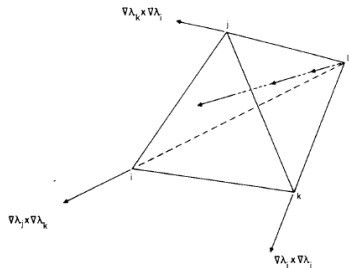
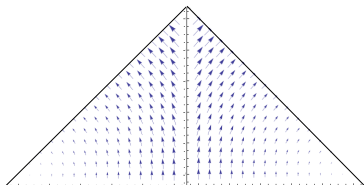
$$\langle 1-x, -y, -z \rangle, \langle x, y, z \rangle$$

$\eta_{ijk} \rightarrow 1$  d.o.f. per face



# Whitney Elements for Primal Meshes

In 3D, it can be shown that the  $\eta_{ij}$  satisfy the  $H(\text{curl})$  constraints and the  $\eta_{ijk}$  satisfy the  $H(\text{div})$  constraints.



See, e.g. Bossavit *Computational Electromagnetism*, 1998.

# Discrete deRham Diagrams

- We now have a basis for finite dimensional subspaces of the deRham Diagram:

$$H^1 \xrightarrow[\text{grad}]{d_0} H(\text{curl}) \xrightarrow[\text{curl}]{d_1} H(\text{div}) \xrightarrow[\text{div}]{d_2} L^2$$

$$\{\lambda_i\} \xrightarrow[(\text{grad})]{\mathbb{D}_0} \{\eta_{ij}\} \xrightarrow[(\text{curl})]{\mathbb{D}_1} \{\eta_{ijk}\} \xrightarrow[(\text{div})]{\mathbb{D}_2} \{\chi_T\}$$

- These are called the **primal cochain spaces** in Discrete Exterior Calculus:

$$\mathcal{C}^0 \xrightarrow[(\text{grad})]{\mathbb{D}_0} \mathcal{C}^1 \xrightarrow[(\text{curl})]{\mathbb{D}_1} \mathcal{C}^2 \xrightarrow[(\text{div})]{\mathbb{D}_2} \mathcal{C}^3$$

- Supposing for a moment we can construct conforming interpolation functions on the dual mesh, we also have a sequence of **dual cochain spaces**:

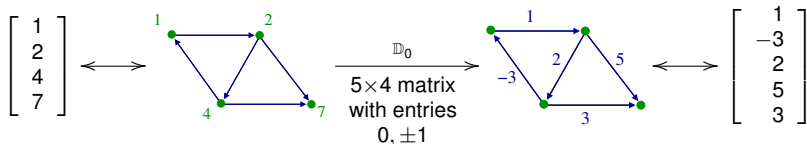
$$\bar{\mathcal{C}}^3 \xleftarrow[(\text{div})]{\mathbb{D}_0^T} \bar{\mathcal{C}}^2 \xleftarrow[(\text{curl})]{\mathbb{D}_1^T} \bar{\mathcal{C}}^1 \xleftarrow[(\text{grad})]{\mathbb{D}_2^T} \bar{\mathcal{C}}^0$$

DESBRUN, HIRANI, LEOK, MARSDEN *Discrete Exterior Calculus*,

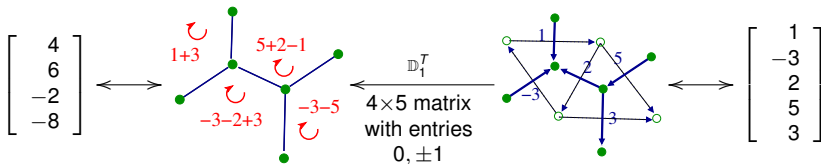
arXiv:math/0508341v2 [math.DG], 2005

# Discrete Exterior Derivative

- The discrete exterior derivative  $\mathbb{D}$  is the transpose of the boundary operator.



- The discrete exterior derivative on the **dual** mesh is  $\mathbb{D}^T$



These cochain vectors and derivative matrices are the building blocks for equation discretization.

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# Discrete Magnetostatics - Primal

Returning to the magnetostatics problem, we can discretize the equations in two ways:

- **Continuous Equations:**

$$\operatorname{div} \mathbf{b} = 0, \quad * \mathbf{b} = \mathbf{h}, \quad \operatorname{curl} \mathbf{h} = \mathbf{j}$$

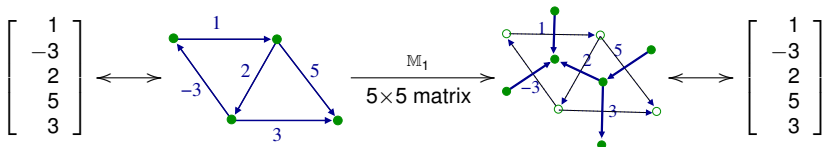
- **'Primal' Discrete Equations**, with  $\mathbf{b}$  as a primal 2-cochain:

$$\mathbb{D}_2 \mathbf{B} = 0, \quad \mathbb{M}_2 \mathbf{B} = \bar{\mathbf{H}}, \quad \mathbb{D}_1^T \bar{\mathbf{H}} = \bar{\mathbf{J}}.$$

- **'Dual' Discrete Equations**, with  $\mathbf{b}$  as a dual 2-cochain:

$$\mathbb{D}_0^T \bar{\mathbf{B}} = 0, \quad \mathbb{M}_1^{-1} \bar{\mathbf{B}} = \mathbf{H}, \quad \mathbb{D}_1 \mathbf{H} = \mathbf{J}.$$

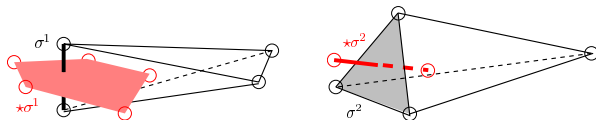
The discrete Hodge Star  $\mathbb{M}$  transfers information between complementary dimensions on **dual** meshes. In this example, we use the identity matrix for  $\mathbb{M}_1$ .



# Discrete Hodge Stars

- Discretization of the Hodge star operator is non-canonical.
- Existing inverse discrete Hodge stars are either too full or too empty for use in discretizations on dual meshes
- We present a novel **dual** discrete Hodge star for this purpose using polyhedral vector interpolation functions

primal mesh simplex  $\sigma^k \iff$  dual mesh cell  $\star\sigma^k$

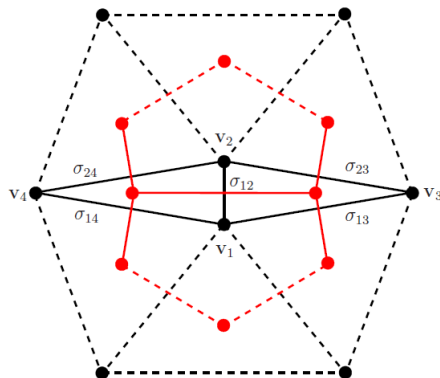


type	reference	definition	$\mathbb{M}_k$	$\mathbb{M}_k^{-1}$
DIAGONAL	[Desbrun et al.]	$(\mathbb{M}_k^{Diag})_{ij} := \frac{ \star\sigma_i^k }{ \sigma_j^k } \delta_{ij}$	diagonal	diagonal
WHITNEY	[Dodziuk],[Bell]	$(\mathbb{M}_k^{Whit})_{ij} := \int_{\mathcal{T}} \eta_{\sigma_i^k} \cdot \eta_{\sigma_j^k}$	sparse	(full)
<b>DUAL</b>	[G, Bajaj]	$((\mathbb{M}_k^{Dual})^{-1})_{ij} := \int_{\mathcal{T}} \eta_{\star\sigma_i^k} \cdot \eta_{\star\sigma_j^k}$	(full)	sparse

# Condition Number of Discrete Hodge Stars

## Theorem [G, Bajaj]

The condition number of  $(\mathbb{M}_k^{Dual})^{-1}$  is governed by different mesh criteria than the condition number of  $\mathbb{M}_k^{Diag}$  and  $\mathbb{M}_k^{Whit}$ .



$$\begin{aligned} \mathbf{v}_1 &= (0, 0) & \mathbf{v}_3 &= (P, \frac{1}{2}) \\ \mathbf{v}_2 &= (0, 1) & \mathbf{v}_4 &= (-P, \frac{1}{2}) \end{aligned}$$

Condition numbers as functions of  $P$ :

$P$	$\mathbb{M}_1^{Diag}$	$\mathbb{M}_1^{Whit}$	$(\mathbb{M}_1^{Dual})^{-1}$
2	6.3	3.2	1.5
5	17.2	9.9	1.3
10	34.6	21.6	1.4
order	$O(P)$	$O(P)$	$O(1)$

*Dual Formulations of Mixed Finite Element Methods. Submitted, 2010.*

# Dual-based Linear Systems

Independence of primal and dual discrete Hodge stars implies **accuracy vs. speed** tradeoffs are possible between primal and dual methods.

**Ex:** Fewer elements in dual mesh  $\rightarrow$  smaller system  $\rightarrow$  faster.

**Ex:** Better condition number in dual system  $\rightarrow$  more accurate.

- **'Primal' Linear System**, with  $b$  as a primal 2-cochain:

$$\begin{aligned} \mathbb{D}_2 \mathbf{B} &= 0, \quad \mathbb{M}_2 \mathbf{B} = \bar{\mathbf{H}}, \quad \mathbb{D}_1^T \bar{\mathbf{H}} = \bar{\mathbf{J}}. \\ \begin{pmatrix} -\mathbb{M}_2 & \mathbb{D}_2^T \\ \mathbb{D}_2 & 0 \end{pmatrix} \begin{pmatrix} \mathbf{B} \\ \bar{\mathbf{P}} \end{pmatrix} &= \begin{pmatrix} -\bar{\mathbf{H}}_0 \\ 0 \end{pmatrix}. \end{aligned}$$

Here,  $\bar{\mathbf{H}}_0 \in \bar{\mathcal{C}}^1$  satisfies  $\mathbb{D}_1^T \bar{\mathbf{H}}_0 = \bar{\mathbf{J}}$  and  $\bar{\mathbf{H}}$  is defined by  $\bar{\mathbf{H}} := \bar{\mathbf{H}}_0 + \mathbb{D}_2^T \bar{\mathbf{P}}$ . Thus  $\mathbb{D}_1^T \bar{\mathbf{H}} = \mathbb{D}_1^T (\bar{\mathbf{H}}_0 + \mathbb{D}_2^T \bar{\mathbf{P}}) = \bar{\mathbf{J}}$  is assured.

- **'Dual' Linear System**, with  $b$  as a dual 2-cochain:

$$\begin{aligned} \mathbb{D}_0^T \bar{\mathbf{B}} &= 0, \quad \mathbb{M}_1^{-1} \bar{\mathbf{B}} = \mathbf{H}, \quad \mathbb{D}_1 \mathbf{H} = \mathbf{J}. \\ \begin{pmatrix} -\mathbb{M}_1^{-1} & \mathbb{D}_0 \\ \mathbb{D}_0^T & 0 \end{pmatrix} \begin{pmatrix} \bar{\mathbf{B}} \\ \mathbf{P} \end{pmatrix} &= \begin{pmatrix} -\mathbf{H}_0 \\ 0 \end{pmatrix}. \end{aligned}$$

Here,  $\mathbf{H}_0 \in \mathcal{C}^1$  satisfies  $\mathbb{D}_1 \mathbf{H}_0 = \mathbf{J}$  and  $\mathbf{H}$  is defined by  $\mathbf{H} := \mathbb{M}_1^{-1} \bar{\mathbf{B}}$ . Thus  $\mathbb{D}_1 \mathbf{H} = \mathbb{D}_1 (\mathbf{H}_0 + \mathbb{D}_0 \mathbf{P}) = \mathbf{J}$  is assured.



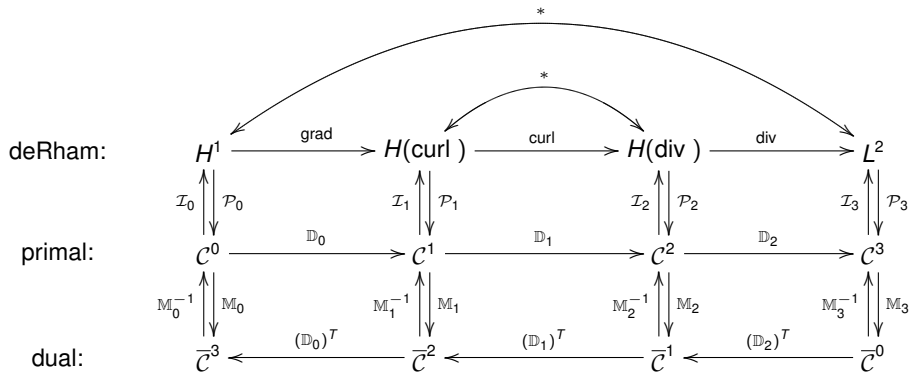
# Dual-based Linear Systems

The duality of the systems is easily visualized via the cochain sequences:

$$\begin{array}{ccccccc}
 \mathcal{C}^0 & \xrightarrow{\text{(grad)}} & \mathcal{C}^1 & \xrightarrow{\text{(curl)}} & \mathcal{C}^2 & \xrightarrow{\text{(div)}} & \mathcal{C}^3 \\
 & & & & & & \\
 & & \textcolor{red}{H} & \xrightarrow{\textcolor{red}{\mathbb{D}_1}} & \textcolor{red}{J} & \textcolor{red}{B} & \xrightarrow{\mathbb{D}_2} 0 \\
 & & \uparrow & & \downarrow & & \\
 & & \textcolor{red}{(M_1)^{-1}} & & \textcolor{red}{M_2} & & \\
 0 & \xleftarrow{\textcolor{red}{(\mathbb{D}_0)^T}} & \textcolor{red}{\bar{B}} & \textcolor{red}{\bar{J}} & \xleftarrow{(\mathbb{D}_1)^T} & \bar{H} & \\
 & & & & & & \\
 \bar{\mathcal{C}}^3 & \xleftarrow{\text{(div)}} & \bar{\mathcal{C}}^2 & \xleftarrow{\text{(curl)}} & \bar{\mathcal{C}}^1 & \xleftarrow{\text{(grad)}} & \bar{\mathcal{C}}^0
 \end{array}$$

# The DEC-deRham Diagram for $\mathbb{R}^3$

We combine the Discrete Exterior Calculus maps with the  $L^2$  deRham sequence.



The combined diagram can be used to formulate dual-based discretizations for many problems including electromagnetics, Darcy flow, and electrodiffusion.

The question remains: How do we construct polyhedral vector interpolation functions?

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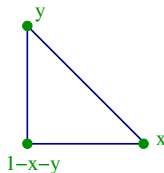
# Scalar Interpolation: Generalized Barycentric Functions

Let  $\Omega$  be a convex polygon in  $\mathbb{R}^2$  with vertices  $\mathbf{v}_1, \dots, \mathbf{v}_n$ . Functions  $\lambda_i : \Omega \rightarrow \mathbb{R}$ ,  $i = 1, \dots, n$  are called **barycentric coordinates** on  $\Omega$  if they satisfy two properties:

- ① **Non-negative:**  $\lambda_i \geq 0$  on  $\Omega$ .
- ② **Linear Completeness:** For any linear function  $L : \Omega \rightarrow \mathbb{R}$ ,  $L = \sum_{i=1}^n L(\mathbf{v}_i) \lambda_i$ .

It can be shown that any set of barycentric coordinates under this definition also satisfy:

- ③ **Partition of unity:**  $\sum_{i=1}^n \lambda_i \equiv 1$ .
- ④ **Linear precision:**  $\sum_{i=1}^n \mathbf{v}_i \lambda_i(\mathbf{x}) = \mathbf{x}$ .
- ⑤ **Interpolation:**  $\lambda_i(\mathbf{v}_j) = \delta_{ij}$ .



## Theorem [Warren, 2003]

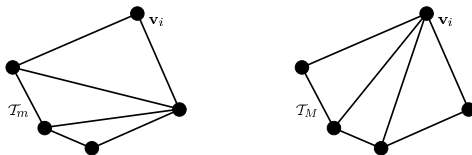
If the  $\lambda_i$  are rational functions of degree  $n - 2$ , then they are unique.

# Triangulation Coordinates

Let  $\mathcal{T}$  be a triangulation of  $\Omega$  formed by adding edges between the  $\mathbf{v}_j$  in some fashion. Define

$$\lambda_{i,\mathcal{T}}^{\text{Tri}} : \Omega \rightarrow \mathbb{R}$$

to be the barycentric function associated to  $\mathbf{v}_i$  on triangles in  $\mathcal{T}$  containing  $\mathbf{v}_i$  and identically 0 otherwise. Trivially, these are barycentric coordinates on  $\Omega$ .



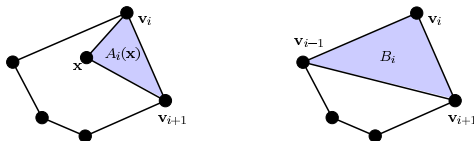
## Theorem [Floater, Hormann, Kós, 2006]

For a fixed  $i$ , let  $\mathcal{T}_m$  denote any triangulation with an edge between  $\mathbf{v}_{i-1}$  and  $\mathbf{v}_{i+1}$ . Let  $\mathcal{T}_M$  denote the triangulation formed by connecting  $\mathbf{v}_i$  to all the other  $\mathbf{v}_j$ . Any barycentric coordinate function  $\lambda_i$  satisfies the bounds

$$0 \leq \lambda_{i,\mathcal{T}_m}^{\text{Tri}}(\mathbf{x}) \leq \lambda_i(\mathbf{x}) \leq \lambda_{i,\mathcal{T}_M}^{\text{Tri}}(\mathbf{x}) \leq 1, \quad \forall \mathbf{x} \in \Omega. \quad (1)$$

# Wachspress Coordinates

Let  $\mathbf{x} \in \Omega$  and define  $A_i(\mathbf{x})$  and  $B_i$  as the areas shown.



Define the Wachspress weight function as

$$w_i^{\text{Wach}}(\mathbf{x}) = B_i \prod_{j \neq i, i-1} A_j(\mathbf{x}).$$

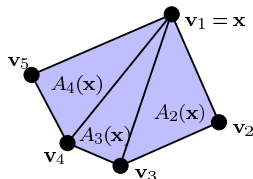
The Wachspress coordinates are then given by the *rational* functions

$$\lambda_i^{\text{Wach}}(\mathbf{x}) = \frac{w_i^{\text{Wach}}(\mathbf{x})}{\sum_{j=1}^n w_j^{\text{Wach}}(\mathbf{x})} \quad (2)$$

# Wachspress Coordinates Example

Let  $\mathbf{x} = \mathbf{v}_1$ .

Note  $A_1(\mathbf{x}) = \text{area of } (\mathbf{v}_1, \mathbf{v}_1, \mathbf{v}_2) = 0$ . Similarly  $A_5(\mathbf{x}) = 0$



$$w_1^{\text{Wach}}(\mathbf{x}) = B_1 A_2(\mathbf{x}) A_3(\mathbf{x}) A_4(\mathbf{x}) = W$$

$$w_2^{\text{Wach}}(\mathbf{x}) = B_2 A_3(\mathbf{x}) A_4(\mathbf{x}) A_5(\mathbf{x}) = 0$$

$$w_3^{\text{Wach}}(\mathbf{x}) = B_3 A_4(\mathbf{x}) A_5(\mathbf{x}) A_1(\mathbf{x}) = 0$$

$$w_4^{\text{Wach}}(\mathbf{x}) = B_4 A_5(\mathbf{x}) A_1(\mathbf{x}) A_2(\mathbf{x}) = 0$$

$$w_5^{\text{Wach}}(\mathbf{x}) = B_5 A_1(\mathbf{x}) A_2(\mathbf{x}) A_3(\mathbf{x}) = 0$$

$$\lambda_1^{\text{Wach}}(\mathbf{x}) = \frac{w_1^{\text{Wach}}(\mathbf{x})}{\sum w_i^{\text{Wach}}(\mathbf{x})} = \frac{W}{W} = 1$$

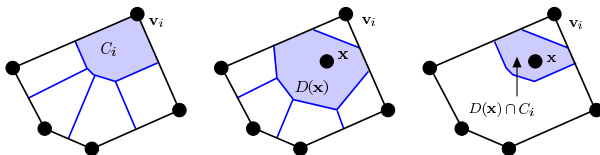
$$\lambda_2^{\text{Wach}}(\mathbf{x}) = \frac{w_2^{\text{Wach}}(\mathbf{x})}{\sum w_i^{\text{Wach}}(\mathbf{x})} = 0$$

Similarly  $\lambda_3^{\text{Wach}}(\mathbf{x}) = \lambda_4^{\text{Wach}}(\mathbf{x}) = \lambda_5^{\text{Wach}}(\mathbf{x}) = 0$ .

This is an illustration of the property  $\lambda_i^{\text{Wach}}(\mathbf{v}_j) = \delta_{ij}$

# Sibson (Natural Neighbor) Coordinates

Let  $P$  denote the set of vertices  $\{\mathbf{v}_i\}$  and define  $P' = P \cup \{\mathbf{x}\}$ .



$$\begin{aligned} C_i &:= |V_P(\mathbf{v}_i)| = |\{\mathbf{y} \in \Omega : |\mathbf{y} - \mathbf{v}_i| < |\mathbf{y} - \mathbf{v}_j|, \forall j \neq i\}| \\ &= \text{area of cell for } \mathbf{v}_i \text{ in Voronoi diagram on the points of } P, \end{aligned}$$

$$\begin{aligned} D(\mathbf{x}) &:= |V_{P'}(\mathbf{x})| = |\{\mathbf{y} \in \Omega : |\mathbf{y} - \mathbf{x}| < |\mathbf{y} - \mathbf{v}_i|, \forall i\}| \\ &= \text{area of cell for } \mathbf{x} \text{ in Voronoi diagram on the points of } P'. \end{aligned}$$

By a slight abuse of notation, we also define

$$D(\mathbf{x}) \cap C_i := |V_{P'}(\mathbf{x}) \cap V_P(\mathbf{v}_i)|.$$

The Sibson coordinates are defined to be

$$\lambda_i^{\text{Sibs}}(\mathbf{x}) := \frac{D(\mathbf{x}) \cap C_i}{D(\mathbf{x})} \quad \text{or, equivalently,} \quad \lambda_i^{\text{Sibs}}(\mathbf{x}) = \frac{D(\mathbf{x}) \cap C_i}{\sum_{j=1}^n D_j(\mathbf{x}) \cap C_j}.$$



# Optimal Coordinates

Let  $g_i : \partial\Omega \rightarrow \mathbb{R}$  be the piecewise linear function satisfying

$$g_i(\mathbf{v}_j) = \delta_{ij}, \quad g_i \text{ linear on each edge of } \Omega.$$

The optimal coordinate function  $\lambda_i^{\text{Opt}}$  is defined to be the solution of Laplace's equations with  $g_i$  as boundary data,

$$\begin{cases} \Delta(\lambda_i^{\text{Opt}}) &= 0, & \text{on } \Omega, \\ \lambda_i^{\text{Opt}} &= g_i. & \text{on } \partial\Omega. \end{cases} \quad (3)$$

These coordinates are optimal in the sense that they minimize the norm of the gradient over all functions satisfying the boundary conditions,

$$\lambda_i^{\text{Opt}} = \operatorname{argmin} \left\{ |\lambda|_{H^1(\Omega)} : \lambda = g_i \text{ on } \partial\Omega \right\}.$$

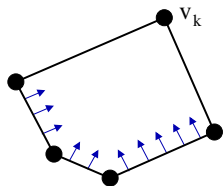
# Polyhedral $H(\text{Curl})$ Vector Interpolation

- Let  $\{\bar{\lambda}_i\}$  denote a set of generalized barycentric coordinates for a polygon (2D) or polyhedra (3D).
- Define for each edge  $\mathbf{v}_i\mathbf{v}_j$ :

$$\bar{\eta}_{ij} := \bar{\lambda}_i \nabla \bar{\lambda}_j - \bar{\lambda}_j \nabla \bar{\lambda}_i$$

## Theorem [G,Bajaj]

Constructing **Whitney-like 1-forms** analogously to the triangular case produces globally  $H(\text{curl})$ -conforming **vector fields**.



PROOF: Consider edge  $\mathbf{v}_i\mathbf{v}_j$  and  $\lambda_k$  associated to a different vertex  $\mathbf{v}_k$ . Then the edge is part of the zero level set of  $\lambda_k$ . Hence  $\nabla \lambda_k$  must be perpendicular to the edge at all points along it and any summand  $\lambda_i \nabla \lambda_k$  has no tangential component on the edge. Therefore, the tangential components only depend on  $\lambda_i$  and  $\lambda_j$ . Hence the  $H(\text{curl})$  conformity constraints are satisfied.  $\square$

To decide which definition of  $\{\bar{\lambda}_i\}$  is suitable, we need error estimates.

# Error Estimates: 2D Scalar Case

The **optimal convergence estimate** for a finite element method bounds the interpolation error in  $H^1$ -norm of an unknown function  $u$  by a constant multiple of the mesh size times the  $H^2$  semi-norm of  $u$ :

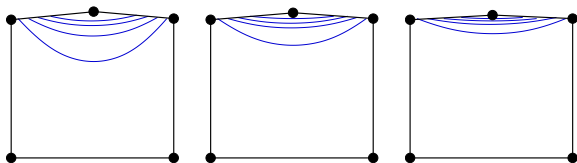
$$\|u - \bar{\mathcal{I}}_0 u\|_{H^1(\Omega)} \leq C \text{diam}(\Omega) |u|_{H^2(\Omega)}, \quad \forall u \in H^2(\Omega). \quad (4)$$

(Note that  $\bar{\mathcal{I}}_0 u$  assumes  $u$  is known or computed at vertices of the dual mesh.)

## Theorem [G, Rand, Bajaj]

Assume certain standard geometric quality conditions on the dual mesh can be guaranteed. Then a **dual formulation** of a finite element method using any of the coordinate systems has the **optimal convergence estimate** on the mesh.

*Error Estimates for Generalized Barycentric Interpolation. Submitted, 2010.*



Example showing necessity of geometric criteria for Wachspress coordinates.

# Future Work

- Efficient computation of  $\bar{\lambda}_i$  basis functions
- Error estimates for polyhedral vector functions
- $H(\text{div})$ -conforming vector elements for polyhedral meshes

# Questions?



- Thank you for inviting me to visit.
- Slides and pre-prints available at <http://www.math.utexas.edu/users/agillette>