

# A Generalization for Stable Mixed Finite Elements

Andrew Gillette

joint work with

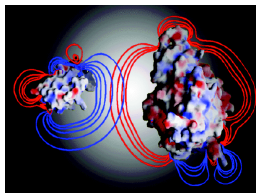
Chandrajit Bajaj

Department of Mathematics  
Institute of Computational Engineering and Sciences  
University of Texas at Austin, USA

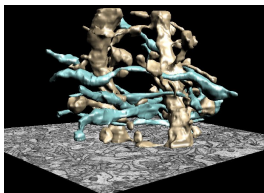
<http://www.math.utexas.edu/users/agillette>

# Motivation

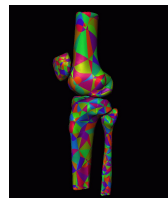
Biological modeling requires **robust** computational methods to solve integral and differential equations over spatially realistic domains.



Electrostatics



Electromagnetics/  
Electrodiffusion



Elasticity

These methods must accommodate

- complicated domain geometry and topology
- multiple variables and operators

# Notions of Robustness

A robust computational method for solving PDEs should exhibit

- **Model Conformity:** Computed solutions are found in a subspace of the solution space for the continuous problem

*Criterion:* Discrete solution spaces replicate the the deRham sequence.

- **Discretization Stability:** The true error between the discrete and continuous solutions is bounded by a multiple of the best approximation error

*Criterion:* The discrete inf-sup condition is satisfied.

- **Bounded Roundoff Error:** Accumulated numerical errors due to machine precision do not compromise the computed solution

*Criterion:* Matrices inverted by the linear solver are well-conditioned.

## Problem Statement

Use the theory of Discrete Exterior Calculus to evaluate the robustness of existing computational methods for PDEs arising in biology and create novel methods with improved robustness.

- 1 Basics of Discrete Exterior Calculus
- 2 Alternative Discretization Pathways
- 3 Stability Criteria from Discrete Hodge Stars

- 1 Basics of Discrete Exterior Calculus
- 2 Alternative Discretization Pathways
- 3 Stability Criteria from Discrete Hodge Stars

# (Continuous) Exterior Calculus

- Differential  $k$ -forms model  $k$ -dimensional physical phenomena.



- The exterior derivative  $d$  generalizes common differential operators.

$$H^1 \xrightarrow[\text{grad}]{d_0} H(\text{curl}) \xrightarrow[\text{curl}]{d_1} H(\text{div}) \xrightarrow[\text{div}]{d_2} L^2$$

- The Hodge Star transfers information between complementary dimensions of primal and dual spaces.

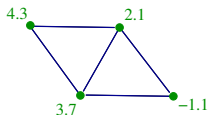
$$\begin{aligned} H^1 &\longleftarrow * \longrightarrow L^2 \\ H(\text{curl}) &\longleftarrow * \longrightarrow H(\text{div}) \end{aligned}$$

## Fundamental “Theorem” of Discrete Exterior Calculus

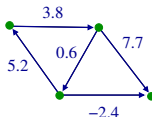
Conforming computational methods must recreate the essential properties of (continuous) exterior calculus on the discrete level.

# Discrete Exterior Calculus

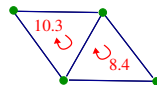
- Discrete differential  $k$ -forms are  $k$ -cochains, i.e. linear functions on  $k$ -simplices.



0-cochain

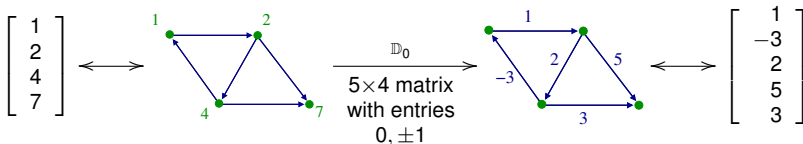


1-cochain



2-cochain

- The discrete exterior derivative  $\mathbb{D}$  is the transpose of the boundary operator.

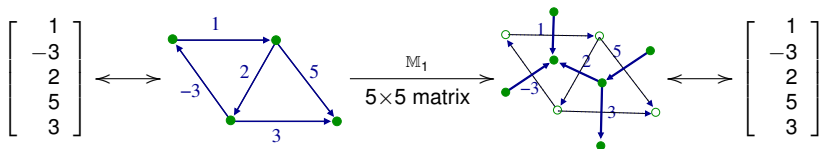


- This creates a discrete analogue of the deRham sequence.

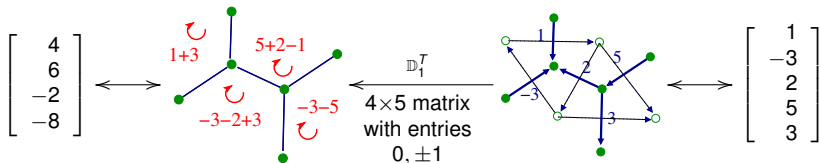
$$\mathcal{C}^0 \xrightarrow[\text{(grad)}]{\mathbb{D}_0} \mathcal{C}^1 \xrightarrow[\text{(curl)}]{\mathbb{D}_1} \mathcal{C}^2 \xrightarrow[\text{(div)}]{\mathbb{D}_2} \mathcal{C}^3$$

# Discrete Exterior Calculus

- The discrete Hodge Star  $\mathbb{M}$  transfers information between complementary dimensions on **dual** meshes. In this example, we use the identity matrix for  $\mathbb{M}_1$ .



- The discrete exterior derivative on the **dual** mesh is  $\mathbb{D}^T$

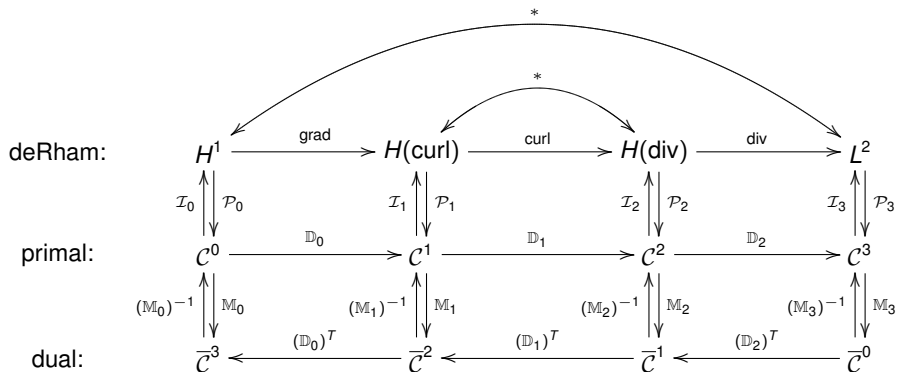


- This creates a **dual** discrete analogue of the deRham sequence.

$$\bar{\mathcal{C}}^3 \xleftarrow[\text{(div)}]{\mathbb{D}_2^T} \bar{\mathcal{C}}^2 \xleftarrow[\text{(curl)}]{\mathbb{D}_1^T} \bar{\mathcal{C}}^1 \xleftarrow[\text{(grad)}]{\mathbb{D}_0^T} \bar{\mathcal{C}}^0$$

# The DEC-deRham Diagram for $\mathbb{R}^3$

We combine the Discrete Exterior Calculus maps with the  $L^2$  deRham sequence.



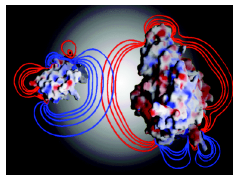
The combined diagram elucidates alternative discretization pathways for finite element methods.

# Outline

- 1 Basics of Discrete Exterior Calculus
- 2 Alternative Discretization Pathways
- 3 Stability Criteria from Discrete Hodge Stars

# Linear Poisson-Boltzmann Equation

$$\operatorname{div}(\epsilon(\vec{x})\nabla\phi(\vec{x})) = \rho_c(\vec{x}) + \bar{\kappa}(\vec{x})\phi(\vec{x}) \text{ in } \mathbb{R}^3$$



$$\begin{aligned}\phi(\vec{x}) &= \text{electrostatic potential} \\ \epsilon(\vec{x}) &= \text{dielectric coefficient} = \begin{cases} \epsilon_I, & \vec{x} \in \Omega \\ \epsilon_E, & \vec{x} \in \mathbb{R}^3 - \Omega \end{cases} \\ \rho_c(\vec{x}) &= \text{charge density from atomic charges} \\ \bar{\kappa}(\vec{x}) &= \text{modified Debye-Huckel parameter}\end{aligned}$$

Exterior calculus formulation:  $d * \epsilon d\phi = f, \quad \phi \in H^1, f \in L^2$

Primal discretization:  $\mathbb{D}_0^T \mathbb{M}_1 \epsilon \mathbb{D}_0 \phi = f$

$$\begin{array}{ccc} \phi & \xrightarrow{\mathbb{D}_0} & \epsilon \mathbb{D}_0 \phi \\ & \downarrow \mathbb{M}_1 & \\ (\mathbb{D}_0)^T \mathbb{M}_1 \epsilon \mathbb{D}_0 \phi & \xleftarrow{(\mathbb{D}_0)^T} & \mathbb{M}_1 \epsilon \mathbb{D}_0 \phi \end{array}$$

Dual discretization:  $\mathbb{D}_2 \epsilon (\mathbb{M}_2)^{-1} (\mathbb{D}_2)^T \phi = f$

$$\begin{array}{ccc} \epsilon (\mathbb{M}_2)^{-1} (\mathbb{D}_2)^T \phi & \xrightarrow{\mathbb{D}_2} & \mathbb{D}_2 \epsilon (\mathbb{M}_2)^{-1} (\mathbb{D}_2)^T \phi \\ \uparrow (\mathbb{M}_2)^{-1} & & \\ (\mathbb{D}_2)^T \phi & \xleftarrow{(\mathbb{D}_2)^T} & \phi \end{array}$$

# DEC for LPBE

The two discretizations inside the DEC-deRham diagram:

$$\begin{array}{ccccccc}
 H^1 & \xrightarrow{\text{grad}} & H(\text{curl}) & \xrightarrow{\text{curl}} & H(\text{div}) & \xrightarrow{\text{div}} & L^2 \\
 \mathcal{I}_0 \updownarrow \mathcal{P}_0 & & \mathcal{I}_1 \updownarrow \mathcal{P}_1 & & \mathcal{I}_2 \updownarrow \mathcal{P}_2 & & \mathcal{I}_3 \updownarrow \mathcal{P}_3 \\
 \phi & \xrightarrow{\mathbb{D}_0} & \epsilon \mathbb{D}_0 \phi & & \epsilon (\mathbb{M}_2)^{-1} (\mathbb{D}_2)^T \phi & \xrightarrow{\mathbb{D}_2} & \mathbb{D}_2 \epsilon (\mathbb{M}_2)^{-1} (\mathbb{D}_2)^T \phi \\
 & & \downarrow \mathbb{M}_1 & & \uparrow (\mathbb{M}_2)^{-1} & & \\
 (\mathbb{D}_0)^T \mathbb{M}_1 \epsilon \mathbb{D}_0 \phi & \xleftarrow{(\mathbb{D}_0)^T} & \mathbb{M}_1 \epsilon \mathbb{D}_0 \phi & & (\mathbb{D}_2)^T \phi & \xleftarrow{(\mathbb{D}_2)^T} & \phi
 \end{array}$$

The definition of the discrete Hodge star matrix  $\mathbb{M}_k$  and its inverse are essential in ensuring the robustness of a primal or dual discretization.

# Further Examples

## Maxwell's Curl Equations

$$\begin{aligned}\nabla \frac{1}{\mu} \times \nabla \times \vec{E} &= \omega^2 \epsilon \vec{E} \\ \nabla \frac{1}{\epsilon} \times \nabla \times \vec{H} &= \omega^2 \vec{H}\end{aligned}$$

$$\begin{array}{ccc} \text{primal:} & \vec{E} & \xrightarrow{\mathbb{D}_1} \mathcal{C}^2 \\ & \updownarrow \scriptstyle \begin{smallmatrix} \mathbb{M}_1 \\ (\mathbb{M}_1)^{-1} \end{smallmatrix} & \updownarrow \scriptstyle \begin{smallmatrix} \mathbb{M}_2 \\ (\mathbb{M}_2)^{-1} \end{smallmatrix} \\ \text{dual:} & \vec{C}^2 & \xleftarrow{(\mathbb{D}_1)^T} \vec{H} \end{array}$$

## Darcy Flow

$$\left\{ \begin{array}{ll} \vec{f} + \frac{k}{\mu} \nabla p &= 0 \quad \text{in } \Omega, \\ \operatorname{div} \vec{f} &= \phi \quad \text{in } \Omega, \\ \vec{f} \cdot \hat{n} &= \psi \quad \text{on } \partial\Omega, \end{array} \right.$$

$$\begin{array}{ccc} \text{primal:} & p & \xrightarrow{\mathbb{D}_0} \begin{matrix} \mathbb{D}_0 p \\ (\mathbb{M}_1^{Dual})^{-1} \vec{f} \end{matrix} \\ & & \uparrow \scriptstyle (\mathbb{M}_1^{Dual})^{-1} \\ \text{dual:} & (\mathbb{D}_0)^T \vec{f} & \xleftarrow{(\mathbb{D}_0)^T} \vec{f} \end{array} \qquad \begin{array}{ccc} & \vec{f} & \xrightarrow{\mathbb{D}_2} \mathbb{D}_2 \vec{f} \\ & \downarrow \scriptstyle \mathbb{M}_2^{Diag} & \\ & \mathbb{M}_2^{Diag} \vec{f} & \xleftarrow{(\mathbb{D}_2)^T} p \end{array}$$

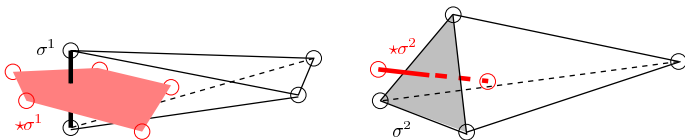
# Outline

- 1 Basics of Discrete Exterior Calculus
- 2 Alternative Discretization Pathways
- 3 Stability Criteria from Discrete Hodge Stars

# Discrete Hodge Star Criteria

A discrete Hodge star transfers information between primal and dual meshes:

**primal mesh simplex**  $\sigma^k \iff$  **dual mesh cell**  $\star\sigma^k$



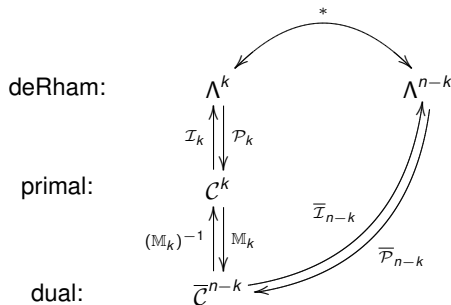
DIAGONAL	[Desbrun et al.]	$(\mathbb{M}_k^{Diag})_{ij} := \frac{ \star \sigma_i^k }{ \sigma_j^k } \delta_{ij}$
WHITNEY	[Dodziuk],[Bell]	$(\mathbb{M}_k^{Whit})_{ij} := \left( \eta_{\sigma_i^k}, \eta_{\sigma_j^k} \right)_{c^k} \quad (\eta_{\sigma^k} = \text{Whitney } k\text{-form for } \sigma^k)$

A **robust** definition of a discrete Hodge star matrix  $\mathbb{M}_k$  and its inverse should provide for:

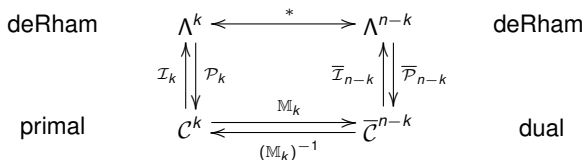
- 1 Commutativity of discrete dual operators
- 2 Local structure of  $\mathbb{M}_k$  and  $\mathbb{M}_k^{-1}$
- 3 Well-conditioned matrices

# Commutativity of Discrete Dual Operators

Given projection to  $(\overline{\mathcal{P}})$  or interpolation from  $(\overline{\mathcal{I}})$  a dual mesh, we have the maps:



Thus, we expect some commutativity of the diagram:



# Commutativity of Discrete Dual Operators

$$\begin{array}{ccc}
 \text{deRham} & \begin{array}{ccc} \Lambda^k & \xleftarrow{*} & \Lambda^{n-k} \\ \mathcal{I}_k \updownarrow \mathcal{P}_k & & \bar{\mathcal{I}}_{n-k} \updownarrow \bar{\mathcal{P}}_{n-k} \\ \mathcal{C}^k & \xrightleftharpoons[\mathbb{M}_k]{(\mathbb{M}_k)^{-1}} & \bar{\mathcal{C}}^{n-k} \end{array} & \text{deRham} \\
 \text{primal} & & \text{dual}
 \end{array}$$

Strong commutativity at  $\mathcal{C}^k$ :

$$*\mathcal{I}_k = \bar{\mathcal{I}}_{n-k}\mathbb{M}_k$$

Weak commutativity at  $\mathcal{C}^k$ :  $\int_{\mathcal{T}} \alpha \wedge *\mathcal{I}_k = \int_{\mathcal{T}} \alpha \wedge \bar{\mathcal{I}}_{n-k}\mathbb{M}_k, \quad \forall \alpha \in \Lambda^k$

**Example:** Discrete Hodge star definitions can be evaluated by this criteria:

$$\text{Using } \mathbb{M}_0^{Diag} : |\mathcal{T}|(\alpha, \lambda_i)_{H^1} = |\star \sigma_i^0| \int_{\star \sigma_i^0} \alpha \mu, \quad \forall \alpha \in H^1$$

$$\text{Using } \mathbb{M}_0^{Whit} : |\mathcal{T}|(\alpha, \lambda_i)_{H^1} = \sum_{\text{vertex } j} (\lambda_i, \lambda_j)_{H^1} \int_{\star \sigma_i^0} \alpha \mu, \quad \forall \alpha \in H^1$$

# Local Structure of $\mathbb{M}_k$ and $\mathbb{M}_k^{-1}$

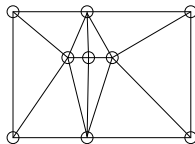
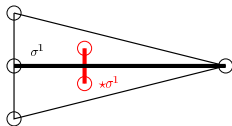
- Existing inverse discrete Hodge stars are either too full or too empty for use in discretizations on dual meshes
- We present a novel **dual** discrete Hodge star for this purpose using generalized barycentric coordinate functions (details in the paper)

type	definition	$\mathbb{M}_k$	$\mathbb{M}_k^{-1}$
DIAGONAL	$(\mathbb{M}_k^{Diag})_{ij} = \frac{ \star \sigma_i^k }{ \sigma_j^k } \delta_{ij}$	diagonal	diagonal
WHITNEY	$(\mathbb{M}_k^{Whit})_{ij} = \left( \eta_{\sigma_i^k}, \eta_{\sigma_j^k} \right)_{C^k}$	sparse	(full)
DUAL	$((\mathbb{M}_k^{Dual})^{-1})_{ij} = \left( \eta_{\star \sigma_i^k}, \eta_{\star \sigma_j^k} \right)_{\bar{C}^k}$	(full)	sparse

# Well-Conditioned Matrices

The condition number of a discrete Hodge star matrix depends on the size of both primal *and* dual mesh elements.

- 1 Primal simplices  $\sigma^k$  satisfy geometric quality measures.
- 2 Dual cells  $\star\sigma^k$  satisfy geometric quality measures.
- 3 The value of  $|\star\sigma^k|/|\sigma^k|$  is bounded above and below.
- 4 The primal and dual meshes do not have large gradation of elements.



We identify which **geometric** criteria are required to provide bounded condition numbers on the various discrete Hodge star matrices.

		<i>i</i>	<i>ii</i>	<i>iii</i>	<i>iv</i>
DIAGONAL	$\mathbb{M}_k^{Diag}$	✓	✓	✓	✓
WHITNEY	$\mathbb{M}_k^{Whit}$	✓			✓
DUAL	$(\mathbb{M}_k^{Dual})^{-1}$		✓		✓

# Questions?



- Slides available at <http://www.ma.utexas.edu/users/agillette>
- This talk was presented at the 2010 Symposium of Solid and Physical Modeling in Haifa, Israel.
- This research was supported in part by NIH contracts R01-EB00487, R01-GM074258, and a grant from the UT-Portugal CoLab project.