# Cauchy-Euler Equations and Method of Frobenius

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Certain singular equations have a solution that is a series expansion. We begin this investigation with **Cauchy-Euler equations**.

## **1** Cauchy-Euler Equations

A second order Cauchy-Euler equation has the form

$$ax^{2}y'' + bxy' + cy = 0 (1)$$

for constants a, b, and c. Thus,  $x_0 = 0$  is a singular point.

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If we try a solution of the form

$$y(x) = x^r,$$

then

$$= ax^{2}y'' + bxy' + cy = (ar(r-1) + br + c)x^{r} = 0.$$

Thus  $x^r$  is a solution if

$$ar(r-1) + br + c = ar_{(b-c)}^{2}r + c = 0$$
<sup>(2)</sup>

The equation (2) is called the **indicial** or **characteristic equation**.

#### Example 1. For

$$x^2y'' - 4xy' + 6y = 0,$$

the indicial equation is

$$0 = r(r-1) - 4r + 6 = r^2 - 5r + 6 = (r-2)(r-3).$$

and

$$y_1(x) = x^2$$
 and  $y_2(x) = x^3$ 

are linearly independent solutions.

As with the constant coefficient equations, we have two additional considerations.

• If the roots in (2) are repeated, i.e., the indicial equation is  $a(r-r_0)^2$ , then the two solutions to (1) are

$$y_1(x) = x^{r_0}$$
 and  $y_2(x) = x^{r_0} \ln x$ .

• If the roots in (2) are complex conjugates,  $\alpha \pm i\beta$ , then the hen the two solutions to (1) are

$$y_1(x) = x^{\alpha} \cos(\beta \ln x)$$
 and  $y_2(x) = x^{\alpha} \sin(\beta \ln x)$ .

**Exercise 2.** • Show that

$$y_{(x)} = c_1 \sqrt{x} + c_2 \sqrt{x} \ln x$$

$$4x'y'' + y = 0$$

• Show that

$$y(x) = \frac{1}{x} \left( c_1 \sin(2\ln x) + c_2 \cos(2\ln x) \right)$$

is a general solution to

$$x^2y'' + 3xy' + 5y = 0$$

This last equation show different behavior possible in Cauchy-Euler equations. The solutions are unbounded and oscillate more and more rapidly near x = 0.

### 2 Method of Frobenius

We moved from second order constant coefficient ordinary differential equations to differential equations having coefficients that are analytic functions of x. The **method of Frobenius** make a similar generalization from the Cauchy-Euler equations.

In this case, we start with

$$a(x)x^{2}y'' + b(x)xy' + c(x)y = 0$$
(3)

and divide so that we have

$$y'' + \frac{b(x)}{xa(x)}y' + \frac{c(x)}{x^2a(x)}y = 0$$
  
$$y'' + p(x)y' + q(x)y = 0$$

So,

$$xp(x) = rac{b(x)}{a(x)}$$
 and  $x^2q(x) = rac{c(x)}{a(x)}$ 

If we have the limits

$$\lim_{x \to 0} xp(x) = \lim_{x \to 0} \frac{b(x)}{a(x)} = p_0 \text{ and } \lim_{x \to 0} x^2 q(x) = \lim_{x \to 0} \frac{c(x)}{a(x)} = q_0$$

Then, for x near zero,

$$p(x) \approx x p_0$$
 and  $q(x) \approx x^2 q_0$ ,

the solutions to (3) should be similar to the Cauchy-Euler equation

$$y'' + \frac{p_0}{x}y' + \frac{q_0}{x^2}y = 0$$
  
$$x^2y'' + xp_0y' + q_0y = 0$$
 (4)

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We turn these observations into a definition.

A singular point  $x_0$  of the differential equation

$$y'' + p(x)y' + q(x)y = 0$$

is said to be a **regular singular point** if both

$$(x - x_0)p(x)$$
 and  $(x - x_0)^2q(x)$ 

are analytic at  $x_0$  Otherwise  $x_0$  is called an **irregular singular point**.

Returning to (4), we again have an indicial equation

$$r(r-1) + p_0 r + q_0 = 0$$

The roots of the indicial equation are called the **exponents** or **indices** of the singularity  $x_0$ .

Example 3. For the Bessel differential equation

$$x^{2}y'' + xy' + (x^{2} - \alpha^{2})y = 0, \quad \alpha \ge 0$$
  
$$y'' + \frac{1}{x}y' + \frac{x^{2} - \alpha^{2}}{x^{2}}y = 0$$
 (5)

To see that  $x_0 = 0$  is a regular singular point, note that

$$\lim_{x \to 0} xp(x) = 1 \quad and \quad \lim_{x \to 0} x^2q(x) = -\alpha^2.$$

The indicial equation is

$$0 = r(r-1) + r - \alpha^2 = r^2 - \alpha^2$$

So the roots are  $\pm \alpha$ 

To begin, let's assume that this equation has distinct real roots,  $r_{-}$  and  $r_{+}$ . The method of Forbenius suggests that we look for solutions of the form

$$y(x) = x^r \sum_{n=0}^{\infty} a_n x^n = \sum_{n=0}^{\infty} a_n x^{n+r}$$

For each of the roots to lead to distinct powers, the difference  $r_{+} - r_{-}$  cannot be an integer.

We apply this to Bessel's equation (5) by first writing series expansions for y, y' and y''.

$$y(x) = \sum_{n=0}^{\infty} a_n x^{n+\alpha} = a_0 x^{\alpha} + a_1 x^{1+\alpha} + a_2 x^{2+\alpha} + a_3 x^{3+\alpha} + a_4 x^{4+\alpha} + \cdots$$

$$y'(x) = \sum_{n=0}^{\infty} (n+\alpha) a_n x^{n+\alpha-1} = \alpha a_0 x^{\alpha-1} + (1+\alpha) a_1 x^{\alpha} + (2+\alpha) a_2 x^{1+\alpha} + (3+\alpha) a_3 x^{2+\alpha} + \cdots$$

$$y''(x) = \sum_{n=0}^{\infty} (n+\alpha) (n+\alpha-1) a_n x^{n+\alpha-2} = \alpha (\alpha-1) a_0 x^{\alpha-2} + (1+\alpha) \alpha a_1 x^{\alpha-1} + (2+\alpha) (1+\alpha) a_2 x^{\alpha} + \cdots$$

For Bessel's equation

$$\begin{aligned} -\alpha^{2}y(x) &= -\sum_{n=0}^{\infty} \alpha^{2}a_{n}x^{n+\alpha} \\ x^{2}y(x) &= \sum_{n=0}^{\infty} a_{n}x^{n+\alpha+2} \\ xy'(x) &= \sum_{n=0}^{\infty} a_{n}x^{n+\alpha+2} \\ xy'(x) &= \sum_{n=0}^{\infty} (n+\alpha)a_{n}x^{n+\alpha} \end{aligned} = -\alpha^{2}a_{0}x^{\alpha} - \alpha^{2}a_{1}x^{1+\alpha} - \alpha^{2}a_{2}x^{2+\alpha} - \alpha^{2}a_{3}x^{3+\alpha} - \alpha^{2}a_{4}x^{4+\alpha} + \cdots \\ = a_{0}x^{2+\alpha} + a_{1}x^{3+\alpha} + a_{2}x^{4+\alpha} + a_{3}x^{5+\alpha} + \cdots \\ = \alpha a_{0}x^{\alpha} + (1+\alpha)a_{1}x^{1+\alpha} + (2+\alpha)a_{2}x^{2+\alpha} + (3+\alpha)a_{3}x^{4+\alpha} + \cdots \end{aligned}$$

$$x^{2}y''(x) = \sum_{n=0}^{\infty} (n+\alpha)(n+\alpha-1)a_{n}x^{n+\alpha} = \alpha(\alpha-1)a_{0}x^{\alpha} + (1+\alpha)\alpha a_{1}x^{\alpha+1} + (2+\alpha)(1+\alpha)a_{2}x^{\alpha+2} + \cdots$$

To line up the indices, we adjust the expression for the  $x^2y(x)$  term.

$$x^{2}y(x) = \sum_{n=2}^{\infty} a_{n-2}x^{n+\alpha} = a_{0}x^{2+\alpha} + a_{1}x^{3+\alpha} + a_{2}x^{4+\alpha} + a_{3}x^{5+\alpha} + \cdots$$

and add. Notice that the sum on the  $x^{\alpha}$  term is

$$-\alpha^2 a_0 + \alpha a_0 + \alpha (\alpha - 1)a_0 = 0$$

. So,  $a_0$  is arbitrary. For the  $x^{\alpha+1}$  term

$$-\alpha^2 a_1 + (1+\alpha)a_1 + (1+\alpha)\alpha a_1 = 0, \quad (2\alpha+1)x_1 = 0, \quad \text{and} \quad a_1 = 0.$$

For the powers of  $x^{n+\alpha}$  we have that

$$0 = (n+\alpha)(n+\alpha-1)a_n + (n+\alpha)a_n - \alpha^2 a_n + a_{n+2}$$
  
=  $((n-\alpha)^2 - \alpha^2 a)_n + a_{n+2}$   
=  $(n+2\alpha)na_n + a_{n-2}$ 

From this we see that the terms  $a_n$  for n odd vanish. Also, notice this recursion relation agrees with the case  $\alpha = 0$  determined earlier.

For the even values of n

$$a_{2} = -\frac{1}{2(2+2\alpha)}a_{0} = -\frac{1}{2(2+2\alpha)}a_{0}$$

$$a_{4} = -\frac{1}{4(4+2\alpha)}a_{2} = \frac{1}{4(4+2\alpha)\cdot2(2+2\alpha)}a_{0} = \frac{1}{2^{4}(2\cdot1)(2+\alpha)(1+\alpha)}a_{0}$$

$$a_{6} = -\frac{1}{6(6+2\alpha)}a_{2} = \frac{1}{6(6+2\alpha)\cdot4(4+2\alpha)\cdot2(2+2\alpha)}a_{0} = \frac{1}{2^{6}(3\cdot)2\cdot1)(3+\alpha)(2+\alpha)(1+\alpha)}a_{0}$$

$$a_{8} = -\frac{2}{8}a_{6} = \frac{1}{4\cdot3\cdot2}a_{0} = -\frac{1}{4!}a_{0}$$

$$\vdots = \vdots = \vdots = \vdots$$

$$a_{2k} = -\frac{1}{k^{2}}a_{2(k-1)} = (-1)^{k}\frac{1}{2^{2k}k!(k+\alpha)_{k}}a_{0}$$

The term  $(a)_k = a(a-1)\cdots(a-k+1)$  is called the falling factorial and is read "x falling k". Thus,  $a_0 J_{\alpha}(x)$  is a solution to (5) where

$$J_{\alpha}(x) = \sum_{n=0}^{\infty} (-1)^k \frac{1}{2^{2k}k!(k+\alpha)_k} x^{2k} = \sum_{n=0}^{\infty} (-1)^k \frac{1}{k!(k+\alpha)_k} \left(\frac{x}{2}\right)^{2k}$$