# Diffusion for a Markov, Divergence-form Generator

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#### **Abstract**

We consider the long-time evolution of solutions to a Schrödinger-type wave equation on a lattice with a Markov random generator. We show that solutions to this problem possess a diffusive scaling limit and compute higher moments.

Based on joint work with Jeffrey Schenker.



## Statement of the Theorem

#### **Theorem**

If  $\psi_t \in \ell^2(\mathbb{Z}^d)$  satisfies

$$\begin{cases}
i\partial_t \psi_t(\mathbf{x}) &= \nabla^{\dagger} \theta_{\omega(t)} \nabla \psi_t(\mathbf{x}) \\
\psi_0(\mathbf{x}) &= \delta_0(\mathbf{x})
\end{cases},$$

then

$$\lim_{\eta \to 0^+} \sum_{\mathbf{x} \in \mathbb{Z}^d} e^{i\sqrt{\eta}k \cdot \mathbf{x}} \mathbb{E}\left( |\psi_{t/\eta}(\mathbf{x})|^2 \right) = e^{-4t \sum_{\theta_1, \theta_2} (k \cdot \theta_1)(k \cdot \theta_2) D_{\theta_1, \theta_2}}.$$



• Consider the standard heat equation

$$\begin{cases}
\partial_t u(x,t) = \Delta u(x,t) & (x,t) \in \mathbb{R}^d \times \mathbb{R}^+ \\
u(x,0) = \delta_0(x) & x \in \mathbb{R}^d
\end{cases}$$

with solution  $u(x, t) = (2\pi t)^{-d/2} e^{-|x|^2/4t}$ .

•  $x \mapsto c_t u(x,t)$  is a p.d.f. on  $\mathbb{R}^d$  with  $c_t = \left(\int_{\mathbb{R}^d} u(x,t) \, dt\right)^{-1}$  the normalizing constant.



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• The p<sup>th</sup> moment of position is given by

$$\int_{\mathbb{R}^d} |x|^p \, c_t u(x,t) \, dx = \frac{c_t \omega_d}{(2\pi t)^{d/2}} \int_0^\infty r^{p+d-1} e^{-\frac{r^2}{4t}} \, dr,$$

where  $\omega_d = |\partial \mathcal{B}(0,1)|$  is the surface area of the unit ball in  $\mathbb{R}^d$ .

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# Diffusive Scaling

#### Definition: Diffusive Scaling

$$\left\{ \begin{array}{ll} t \mapsto \frac{1}{\eta} t \\ x \mapsto \frac{1}{\sqrt{\eta}} x \end{array} \right. \text{ as } \eta \to 0^+$$

- Question: The problem under consideration is defined on the lattice  $\mathbb{Z}^d$ . How do we scale a discrete space?
- Answer: Mollify.



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- $h \in C_c^{\infty}(\mathbb{R}^d)$ ,  $\int h \, dx = 1$ ,  $h \geq 0$
- Under diffusive scaling, if the convolution  $h*|\psi_t|^2$  converges (weakly) to a solution of the heat equation, then we say that the model exhibits *diffusion*.
- A Fourier transform removes the mollifier from our diffusion criterion.
- Diffusion Criterion:

$$\sum_{\mathbf{x}\in\mathbb{Z}^d} e^{i\sqrt{\eta}k\cdot\mathbf{x}} |\psi_{t/\eta}(\mathbf{x})|^2 \to e^{-Dt|\mathbf{k}|^2}, \quad \mathbf{k}\in\mathbb{T}^d$$



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$$\sum_{\mathbf{x} \in \mathbb{Z}^d} e^{i\sqrt{\eta}k \cdot \mathbf{x}} \mathbb{E}(|\psi_{t/\eta}(\mathbf{x})|^2)$$

$$= -\frac{1}{2\pi i} \int_{\Gamma} e^{-t\mathbf{z}} \left\langle \delta_0 \otimes \mathbf{1}, \frac{\eta}{i\hat{L}_{\sqrt{\eta}k} + B - \eta \mathbf{z}} \delta_0 \otimes \mathbf{1} \right\rangle d\mathbf{z}$$

- Notes:
  - The LHS is (almost) the diffusion criterion.
  - The expectation allows us to use a Feynman-Kac-Pillet formula.
  - FKP allows us to express the expectation as a matrix element of the semigroup  $e^{-t(i\hat{L}\sqrt{\eta}k+B)}$ ,
  - which can be understood by the holomorphic functional calculus:

$$e^{t(i\hat{L}_{\sqrt{\eta}k}+B)}=rac{1}{2\pi i}\int_{\Gamma}e^{tZ}rac{1}{i\hat{L}_{\sqrt{\eta}k}+B-z}\,\mathrm{d}z$$



### • Key step:

$$\begin{split} & \sum_{\mathbf{x} \in \mathbb{Z}^d} e^{i\sqrt{\eta}k \cdot \mathbf{x}} \mathbb{E}(|\psi_{t/\eta}(\mathbf{x})|^2) \\ & = -\frac{1}{2\pi i} \int_{\Gamma} e^{-t\mathbf{z}} \left\langle \delta_0 \otimes \mathbf{1}, \frac{\eta}{i\hat{L}_{\sqrt{\eta}k} + B - \eta \mathbf{z}} \delta_0 \otimes \mathbf{1} \right\rangle \, d\mathbf{z} \end{split}$$

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• We have reduced the problem to understanding:

$$\lim_{\eta \to 0^+} \left\langle \delta_0 \otimes 1, \frac{\eta}{i\hat{L}_{\sqrt{\eta}k} + B - \eta z} \, \delta_0 \otimes 1 \right\rangle.$$

- From here,
  - use projections and the Schur complement formula.
  - construct a symmetric operator  $D_k$ , which is a lower bound for the matrix element in question. Use this to show the limit exists and is of the desired form.
- Higher Moments?

$$\lim_{\eta \to 0^+} \sum_{\mathbf{x} \in \mathbb{Z}^d} e^{i\sqrt{\eta}k \cdot \mathbf{x}} \mathbb{E}\left(|\psi_{t/\eta}(\mathbf{x})|^2\right) = e^{-4t\sum_{e_1,e_2}(k \cdot e_1)(k \cdot e_2)D_{e_1,e_2}},$$



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## Thank You & References

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